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# Sample-based Hamiltonian simulation via density matrix exponentiation

- This is an approach for performing Hamiltonian simulation by interacting a system of interest w/ program states, according to a repeated, fixed interaction.
- Before getting into it, we should review the notion of density matrix and partial trace.
- So far, we have described q. states as state vectors, Another way of doing so is  $\beta$

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Via density matrices.

- Density matrix corresponding to state vector  $|\psi\rangle$  is outer product

$$|\psi\rangle\langle\psi|$$

- Notice it has the benefit of removing a global phase because D.M. for  $|\psi\rangle$  is

$$|\psi\rangle\langle\psi| \quad \& \quad \text{that for}$$

$$e^{i\phi}|\psi\rangle \quad \text{is}$$

$$e^{i\phi}|\psi\rangle\langle\psi|e^{-i\phi} = |\psi\rangle\langle\psi|$$

- If we are unsure about precise state of a system, and instead only know that it is described by

probabilistic ensemble  $\{p(x), |\psi_x\rangle\}_{x \in X}$

then corresponding density matrix is

$$\sum_x p(x) |\psi_x\rangle\langle\psi_x| = \rho$$

If we perform a measurement  $\{\Pi_y\}_{y \in Y}$ ,  
outcome probability is

$$p(y) = \text{Tr}[\Pi_y \rho]$$

$$= \sum_x p(x) \text{Tr}[\Pi_y |\psi_x\rangle\langle\psi_x|]$$

$$= \sum_x p(x) \langle\psi_x | \Pi_y | \psi_x\rangle$$

$$= \sum_x p(x) \underbrace{\|\Pi_y |\psi_x\rangle\|_2^2}$$

$$\sum_x p(x) P(y|x)$$

consistency w/ law of total probability

Also, ~~with~~

$$|4\rangle \rightarrow U|4\rangle \Leftrightarrow |4\rangle\langle 4| \rightarrow U|4\rangle\langle 4|U^\dagger \quad (4)$$

other system we need  $\rho_B$   
partial trace.

State of two-party ~~Herbert~~  
system

$$\rho_{AB}$$

reduced density matrix

$$\rho_A = \text{Tr}_B[\rho_{AB}]$$

state

$$\rho_A = \sum_i (I_A \otimes \langle i|_B) \rho_{AB} (I_A \otimes |i\rangle_B)$$

of Alice's  
system,  
used to predict  
outcomes.

of measurements  
on Alice's  
system alone

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von-Neumann equation is a  
consequence of Schrödinger equation

Recall that

$$i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle = H |\psi(t)\rangle$$

(time-independent Schrö.  
equation)

Consider evolution of  
density matrix

$$i\hbar \frac{\partial}{\partial t} (|\psi(t)\rangle \langle \psi(t)|)$$

$$= i\hbar \left( \left| \frac{\partial}{\partial t} \psi(t) \right\rangle \langle \psi(t)| + |\psi(t)\rangle \left\langle \frac{\partial}{\partial t} \psi(t) \right| \right)$$

$$= \left( i\hbar \left| \frac{\partial}{\partial t} \psi(t) \right\rangle \right) \langle \psi(t)| - |\psi(t)\rangle \left( \left\langle \frac{\partial}{\partial t} \psi(t) \right| - i\hbar \right)$$

product rule

$$= H |\psi(t)\rangle \langle \psi(t)| - |\psi(t)\rangle \langle \psi(t)| H$$

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ble conj. transpose of Schröd. eqn.

is

$$-i\hbar \left\langle \frac{\partial}{\partial t} \psi(t) \right| = \langle \psi(t) | H$$

↑  
Hermitian

$$\Rightarrow i\hbar \frac{\partial}{\partial t} [ \langle \psi(t) | \langle \psi(t) | ] \\ = [ H, \langle \psi(t) | \langle \psi(t) | ]$$

Now consider general

$$\rho(t) = \sum_x P_x(t) | \psi_x(t) \rangle \langle \psi_x(t) |$$

$$i\hbar \frac{\partial}{\partial t} \rho(t)$$

$$= i\hbar \left( \sum_x \left[ \frac{\partial}{\partial t} P_x(t) \right] | \psi_x(t) \rangle \langle \psi_x(t) | \right. \\ \left. + \sum_x P_x(t) \frac{\partial}{\partial t} [ | \psi_x(t) \rangle \langle \psi_x(t) | ] \right)$$

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$$= i\hbar \sum_x \left[ \frac{\partial}{\partial t} P_x(t) \right] |\psi_x(t)\rangle \langle \psi_x(t)|$$

$$+ \sum_x P_x(t) [H, |\psi_x(t)\rangle \langle \psi_x(t)|]$$

$$= i\hbar \sum_x \left[ \frac{\partial}{\partial t} P_x(t) \right] |\psi_x(t)\rangle \langle \psi_x(t)|$$

$$+ [H, \rho(t)]$$

for a closed system,

probabilities do not change

$$\text{w/ time} \Rightarrow \frac{\partial}{\partial t} P_x(t) = 0$$

$$\Rightarrow i\hbar \frac{\partial}{\partial t} \rho(t) = [H, \rho(t)]$$

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Equivalent way to think about this:

Initial state is

$$\rho(0) = \sum_x p_x(0) |\psi_x(0)\rangle \langle \psi_x(0)|$$

solution of Schröd. eqn.

is

$$|\psi(t)\rangle = e^{-iHt/\hbar} |\psi(0)\rangle$$

⇒ each pure state evolves according to

$$|\psi_x(t)\rangle = e^{-iHt/\hbar} |\psi_x(0)\rangle$$

⇒ density matrix @ some  $t$  is

$$\rho(t) = \sum_x p_x(0) e^{-iHt/\hbar} |\psi_x(0)\rangle \langle \psi_x(0)|$$

$$= e^{-iHt/\hbar} \rho(0) e^{iHt/\hbar}$$

† is solution of von Neumann eqn.



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Idea of density matrix exponentiation:

— assume ~~the~~ Hamiltonian  $H$  is positive semidefinite & has trace one, so that it can be encoded as a density matrix, call it  $\sigma$ .

— Then suppose we have sample access to  $\sigma$ , such that we can call it multiple times

then perform an interaction such that

$$\rho \otimes \sigma \otimes \dots \otimes \sigma \rightarrow \underbrace{\rho}_{\mathcal{N}} \otimes e^{-i\sigma t} \otimes \underbrace{\rho}_{\mathcal{N}} e^{i\sigma t}$$

Consider that desired evolution  
for short time  $\Delta$  is

$$\begin{aligned} \rho(\Delta) &= e^{-i\sigma\Delta} \rho e^{i\sigma\Delta} \\ &= (\mathbb{I} + (-i\sigma\Delta) + o(\Delta^2)) \rho (\mathbb{I} + i\sigma\Delta + o(\Delta^2)) \\ &= \rho - i\sigma\rho\Delta + i\rho\sigma\Delta + o(\Delta^2) \\ &= \rho - i[\sigma, \rho]\Delta + o(\Delta^2) \end{aligned}$$

Observe that

$$\frac{\rho(\Delta) - \rho}{\Delta} = -i[\sigma, \rho] + o(\Delta)$$

$\lim_{\Delta \rightarrow 0}$  gives von Neumann equation

$$\dot{\rho} = -i[\sigma, \rho]$$

so this kind of short time step  
expansion is similar to how  
differential equations are solved

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numerically, except we're considering  
this on a q. computer.

Now consider this

take initial state to

be  $\rho_{00}$  &

Hamiltonian to be  $\text{SWAP} = F$

then

$$\begin{aligned} & e^{-iF\Delta} (\rho_{00}) e^{iF\Delta} \\ &= (\mathbb{I} + (-iF\Delta) + o(\Delta^2)) (\rho_{00}) (\mathbb{I} + iF\Delta + o(\Delta^2)) \\ &= \rho_{00} - iF(\rho_{00})\Delta + i(\rho_{00})F\Delta + o(\Delta^2) \end{aligned}$$

Now suppose we take partial trace  
over system 2

This implies that

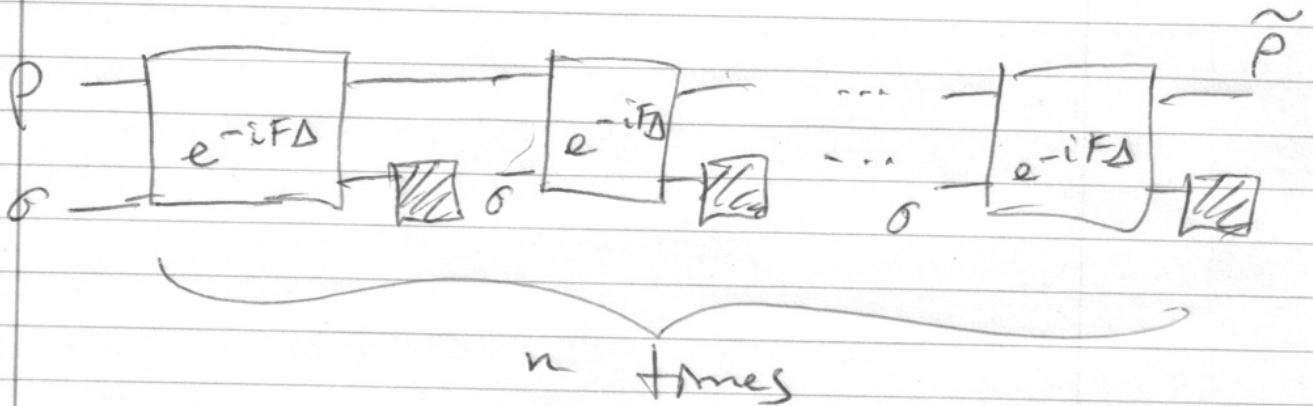
$$\begin{aligned}
 & \text{Tr}_2 \left[ e^{-iF\Delta} (\rho \otimes \sigma) e^{iF\Delta} \right] \\
 &= \text{Tr}_2 \left[ \rho \otimes \sigma \right] - i \text{Tr}_2 \left[ F(\rho \otimes \sigma) \right] \Delta \\
 &\quad + i \text{Tr}_2 \left[ (\rho \otimes \sigma) F \right] \Delta + o(\Delta^2) \\
 &= \rho \otimes \sigma - i \sigma \rho \Delta + i \rho \sigma \Delta + o(\Delta^2) \\
 &= \rho - i [\sigma, \rho] \Delta + o(\Delta^2)
 \end{aligned}$$

⇒ SWAP interaction for short  
time & program state  
can simulate dynamics  
correctly for short time  
 $\Delta$ .

Idea then is to pick

$$\Delta = \frac{t}{n} \quad \& \text{ execute the}$$

following protocol:



& we can guarantee that

$$\tilde{\rho} \approx_{\epsilon} e^{-i\sigma t} \rho e^{i\sigma t}$$

$$\text{if } n = \frac{t^2}{\epsilon}$$

Proper error analysis:

Consider that

$$\begin{aligned} & \text{Tr}_2 [ e^{-iF\Delta} (\rho_{\sigma\sigma}) e^{iF\Delta} ] \\ &= \text{Tr}_2 \left[ \left( I - iF\Delta + \frac{I\Delta^2}{2} + o(\Delta^3) \right) \right. \\ & \quad \left. \left( I + iF\Delta - \frac{I\Delta^2}{2} + o(\Delta^3) \right) \right] \end{aligned}$$

d/c  $(-iF\Delta)^2 = -I\Delta^2$

$$\begin{aligned} &= \text{Tr}_2 \left[ \rho_{\sigma\sigma} - iF(\rho_{\sigma\sigma})\Delta + i(\rho_{\sigma\sigma})F\Delta \right. \\ & \quad \left. + \underbrace{F(\rho_{\sigma\sigma})F}_{\sigma\rho} \Delta^2 - \rho_{\sigma\sigma} \Delta^2 + o(\Delta^3) \right] \\ &= \rho - i[\sigma, \rho]\Delta + (\sigma - \rho)\Delta^2 + o(\Delta^3) \end{aligned}$$

⇒ State @ kth step of algorithm is

$$\begin{aligned} \tilde{\rho}^{(k)} &= \tilde{\rho}^{(k-1)} - i[\sigma, \tilde{\rho}^{(k-1)}]\Delta \\ & \quad + (\sigma - \tilde{\rho}^{(k-1)})\Delta^2 + o(\Delta^3) \end{aligned}$$

$$\begin{aligned} \downarrow \tilde{\rho}^{(k-1)} &= \tilde{\rho}^{(k-2)} - i[\sigma, \tilde{\rho}^{(k-2)}] \Delta \\ &\quad + (\sigma - \tilde{\rho}^{(k-2)}) \Delta^2 + o(\Delta^3) \end{aligned}$$

substitute  $\tilde{\rho}^{(k-1)}$  expression into  
1st line

$\downarrow$  get

$$\begin{aligned} \tilde{\rho}^{(k)} &= \tilde{\rho}^{(k-2)} - i[\sigma, \tilde{\rho}^{(k-2)}] \Delta \\ &\quad + (\sigma - \tilde{\rho}^{(k-2)}) \Delta^2 + o(\Delta^3) \\ &\quad - i[\sigma, \tilde{\rho}^{(k-2)}] \Delta \\ &\quad - i[\sigma, -i[\sigma, \tilde{\rho}^{(k-2)}]] \Delta^2 \\ &\quad + (\sigma - \tilde{\rho}^{(k-2)}) \Delta^2 + o(\Delta^3) \end{aligned}$$

$$\begin{aligned} &= \tilde{\rho}^{(k-2)} - i[\sigma, \tilde{\rho}^{(k-2)}] 2\Delta \\ &\quad + (\sigma - \tilde{\rho}^{(k-2)}) 2\Delta^2 \\ &\quad - i[\sigma, -i[\sigma, \tilde{\rho}^{(k-2)}]] \Delta^2 \end{aligned}$$

Repeating this recursion ~~steps~~

$n$  times starting from initial state gives

$$\tilde{\rho}^{(n)} = \tilde{\rho}^{(0)} - i [\sigma, \tilde{\rho}^{(0)}] n \Delta$$

$$+ (\sigma - \tilde{\rho}^{(0)}) n \Delta^2$$

$$- i [\sigma, -i [\sigma, \tilde{\rho}^{(0)}]] (1+2+\dots+n-1) \Delta^2$$

$$- [\sigma, [\sigma, \tilde{\rho}^{(0)}]] + o(\Delta^3)$$

$$1+2+\dots+n-1 = \frac{n(n-1)}{2}$$

Ideal evolution is  $\rho$  for time  $t = n\Delta$

$$e^{-i\sigma n\Delta} \rho e^{i\sigma n\Delta}$$

$$= \left( \mathbb{I} - i\sigma n\Delta - \frac{\sigma^2 n^2 \Delta^2}{2} \right) \rho \left( \mathbb{I} + i\sigma n\Delta - \frac{\sigma^2 n^2 \Delta^2}{2} + o(\Delta^3) \right)$$



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$$= \rho - i[\sigma, \rho]n\Delta - \frac{1}{2}[\sigma, [\sigma, \rho]]n^2\Delta^2 + o(\Delta^3)$$

then subtracting ideal  
from actual gives

$$(\sigma - \hat{\rho}^{(0)})n\Delta^2 - [\sigma, [\sigma, \hat{\rho}^{(0)}]]n\Delta^2$$

$$= O(n\Delta^2) = O\left(\frac{t^2}{n}\right) \quad \Delta = t/n$$

$\Rightarrow$  if desired error is  $\epsilon$

then pick  $\epsilon = \frac{t^2}{n}$

$$\Rightarrow n = \frac{t^2}{\epsilon}$$

$$\Rightarrow n = O\left(\frac{t^2}{\epsilon}\right)$$