

## Lecture 22

①

orthogonality of Bessel functions

Consider  $\sin x$  &  $\cos x$ .

these are similar to  $J_p(x)$  &  $Y_p(x)$ .

$$\sin x = 0 \Leftrightarrow x = n\pi \text{ for } n \in \mathbb{Z}$$

same as

$$\sin(n\pi x) \Leftrightarrow x \in \mathbb{Z}$$

$$\sin''(n\pi x) = -(n\pi)^2 \sin(n\pi x)$$

$$\Rightarrow \sin''(n\pi x) + (n\pi)^2 \sin(n\pi x) = 0$$

Now consider  $J_p(x)$  &  $Y_p(x)$

As stated above, these are like

$\sin x$  &  $\cos x$ , respectively

(2)

The roots of  $J_p(x)$  are not integers but we can enumerate them as  $J_p(\alpha_p) = 0, J_p(\beta_p) = 0, \dots$

where  $\alpha_p, \beta_p, \gamma_p$  are the roots of  $J_p(x)$

$$\Rightarrow \text{ @ } x=1 \quad J_p(\alpha_p x) = 0, \\ J_p(\beta_p x) = 0, \dots$$

Recall that  $\{\sin(n\pi x)\}_n$  form an orthonormal set over  $[0, 1]$ .

$$\text{that is, } \int_0^1 \sin(n\pi x) \sin(m\pi x) dx = 0 \\ \text{if } n \neq m$$

$$\int_0^1 \sin(n\pi x) \sin(m\pi x) dx = 1/2 \\ \text{if } n = m$$

(3)

Our goal is to prove that

$$\int_0^1 x J_p(\alpha_p x) J_p(\beta_p x) dx = 0$$

for  
 $\alpha_p \neq \beta_p$

Consider that

$$x^2 J_p''(x) + x J_p'(x) + (x^2 - p^2) J_p(x) = 0$$

~~can be scale by roots  $\alpha_p$  or  $\beta_p$~~   
as  $\frac{d}{dx} J_p(x)$

Now substitute  $z = \alpha x$  to get

$$z^2 \frac{d^2}{dz^2} J_p(z) + z \frac{d}{dz} J_p(z) +$$

$$(z^2 - p^2) J_p(z) = 0$$

use chain rule to get

$$\frac{dJ_p(z)}{dz} \frac{dx}{dz} \frac{d}{dx} J_p(\alpha x) = \frac{1}{\alpha} \frac{d}{dx} J_p(\alpha x)$$

(4)

$$\begin{aligned} & (\alpha^2)^2 \frac{1}{\alpha^2} \frac{d^2}{dx^2} J_p(\alpha x) + \alpha x \frac{1}{\alpha} \frac{d}{dx} J_p(\alpha x) \\ & + (\alpha^2 x^2 - p^2) J_p(\alpha x) = 0 \end{aligned}$$

$$\Rightarrow x^2 \frac{d^2}{dx^2} J_p(\alpha x) + x \frac{d}{dx} J_p(\alpha x) + (\alpha^2 x^2 - p^2) J_p(\alpha x) = 0$$

Now use  $x \frac{d}{dx} [x \frac{d}{dx} y] = x \frac{dy}{dx} + x^2 \frac{d^2 y}{dx^2}$

to rewrite as

$$x \frac{d}{dx} \left[ x \cdot \frac{d}{dx} J_p(\alpha x) \right] + (\alpha^2 x^2 - p^2) J_p(\alpha x) = 0$$

same for  $\beta$

$$x \frac{d}{dx} \left[ x \frac{d}{dx} J_p(\beta x) \right] + (\beta^2 x^2 - p^2) J_p(\beta x) = 0$$

5

Recall that  $\alpha$  &  $\beta$  are arbitrary roots of  $J_p(x)$ .

$$u(x) \equiv J_p(\alpha x)$$

$$v(x) \equiv J_p(\beta x)$$

to simplify notation.

Goal is to prove that

$$\int_0^1 x J_p(\alpha x) J_p(\beta x) dx = \frac{1}{2} [J_p'(\alpha)]^2 \delta_{\alpha, \beta}$$

Proof: Suppose  $\alpha \neq \beta$ , where  $\delta_{\alpha, \beta} = \begin{cases} 0 & \text{if } \alpha \neq \beta \\ 1 & \text{if } \alpha = \beta \end{cases}$   
we have that

$$x [x u']' + (\alpha^2 x^2 - p^2) u = 0$$

$$x [x v']' + (\beta^2 x^2 - p^2) v = 0$$

Multiply 1st by  $v$ , 2nd by  $u$   
subtract & divide by  $x$  to get

$$v(xu')' - u(xv')' + (\alpha^2 - \beta^2)xuv = 0 \quad (6)$$

can rewrite as

$$\frac{d}{dx} [vxu' - uxv'] + (\alpha^2 - \beta^2)xuv = 0$$

Now integrate from 0 to 1 to  
get

$$\underbrace{(vxu' - uxv')}'_0 + (\alpha^2 - \beta^2) \int_0^1 xuv \, dx = 0$$

lower limit = 0 b/c  $x$  is present

$\&$   $v, u', u, v'$   
are finite

upper limit = recall that

$$u = J_p(\alpha) = 0 \quad \& \quad v = J_p(\beta) = 0$$

since  $\alpha$   $\&$   $\beta$  are zeros of  
 $J_p(x)$ .

(7)

Then we get

$$(\alpha^2 - \beta^2) \int_0^1 x u v dx = 0$$

or

$$(\alpha^2 - \beta^2) \int_0^1 x J_p(\alpha x) J_p(\beta x) dx = 0$$

$$\Rightarrow \int_0^1 x J_p(\alpha x) J_p(\beta x) dx = 0$$

if  $\alpha \neq \beta$ .

To see the result for  $\alpha = \beta$ ,

we employ a limiting argument.

Let  $\alpha$  be a root of  $J_p(x)$

& let  $\beta \neq \alpha$  w/  $\beta$  not a root.

- Note that

$$(v x u' - u x v') \Big|_0^1 + (\alpha^2 - \beta^2) \int_0^1 x u v dx = 0$$

still holds but  $J_p(\beta) \neq 0$  &  $J_p(\alpha) = 0$ .

8

rewrite as

$$\left[ J_p(\beta x) \times J_p'(\alpha x) - J_p(\alpha x) \times J_p'(\beta x) \right] \Big|_0^1 \\ + (\alpha^2 - \beta^2) \int_0^1 x J_p(\alpha x) J_p(\beta x) dx = 0$$

lower limit of 1st term = 0

$$\text{upper limit is } = J_p(\beta) J_p'(\alpha) \cdot \alpha$$

b/c  $J_p(\alpha \cdot 1) = 0$ , so right term

goes away

$$\nabla \lim_{x \rightarrow 1} \frac{d}{dx} J_p(\alpha x)$$

$$= \lim_{x \rightarrow 1} \frac{dz}{dx} \cdot \frac{d}{dz} J_p(z) \quad \text{Set } z = \alpha x$$

$$= \alpha \cdot \lim_{x \rightarrow 1} \frac{d}{dz} J_p(z)$$

$$= \alpha \cdot J_p'(\alpha)$$



(9)

$$\Rightarrow \frac{J_p(\beta) \cdot \alpha \cdot J_p'(\alpha)}{\beta^2 - \alpha^2} = \int_0^1 x J_p(\alpha x) J_p(\beta x) dx$$

Now take  $\lim_{\beta \rightarrow \alpha}$ . Both top + bottom

of LHS  $\rightarrow 0$ . Apply

L'Hospital's rule + differentiate

top + bottom wrt  $\beta$ :

$$\lim_{\beta \rightarrow \alpha} \frac{J_p'(\beta) \cdot \alpha \cdot J_p'(\alpha)}{2\beta} = \frac{1}{2} [J_p'(\alpha)]^2$$

So we have proven that

$$\int_0^1 x J_p(\alpha x) J_p(\beta x) dx = \frac{1}{2} [J_p'(\alpha)]^2 \delta_{\alpha\beta}$$

where  $\alpha$  +  $\beta$  are roots of  $J_p(x)$ .

(10)

2 ways to look @ this

1) The functions

$$\tilde{j}_p(\alpha x) \equiv \sqrt{2} [J_p'(\alpha)]^{-1} \cdot \sqrt{x} J_p(\alpha x)$$

where  $\alpha$   
is a root

are orthonormal on  $x \in [0, 1]$ .

That is

$$\int_0^1 \tilde{j}_p(\alpha_p x) \tilde{j}_p(\beta_p x) dx = \delta_{\alpha, \beta}$$

As such, we can expand  
well-behaved functions in terms  
of Bessel's functions:

$$f(x) = \sum_{\alpha} c_{\alpha} \tilde{j}_p(\alpha x)$$

where

$$c_{\alpha} = \int_0^1 \tilde{j}_p(\alpha x) f(x) dx.$$

(11)

The functions

$$q_p(\alpha_p(x)) \equiv \sqrt{2} [J_p'(\alpha)]^{-1} J_p(\alpha x)$$

are orthonormal w/ respect to

weight function  $\rho(x) = x$ , s.t.,

$$\int_0^1 \rho(x) q_p(\alpha_p(x)) q_p(\beta_p(x)) dx = \delta_{\alpha, \beta}$$

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Now show page 604

Table of Asymptotic forms

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Fuchs's theorem:

Now we've discussed two examples of differential equations, Legendre & Bessel, solvable by the Frobenius method

(12)

of generalized power series.

- There are other examples of these kinds of functions:

- they are the solutions of parametrized diff. eq's.

- have orthogonality properties & other functions can be expanded in terms of them.

- generating functions for them.

- physical problems whose solutions involve them.

- Might wonder if all diff. eq's might be solved in this way.

- General theorem due to Fuchs states when this is possible:

Fuchs's theorem for 2<sup>nd</sup> order diff-eqs is

$$y'' + f(x)y' + g(x)y = 0$$

If  $xf(x)$  &  $x^2g(x)$  are  $\sum_{n=0}^{\infty} a_n x^n$  expandable in convergent power series, then diff. eq. is regular at the origin.

Theorem of Fuchs says that these conditions are necessary & sufficient for solution to consist of

1) two Frobenius series, i.e.,

$$y(x) = x^s \sum_{n=0}^{\infty} c_n x^n + x^t \sum_{n=0}^{\infty} d_n x^n$$

or

2) one solution being a Frobenius series  $S_1(x)$  & a second solution  $S_1(x) \ln x + S_2(x)$  w/  $S_2(x)$  Frobenius.