

Lecture 17

26 MAR 2014

On the optimality of Grover's algorithm

We can prove that $\Omega(\sqrt{2^n})$

queries are needed for any

quantum algorithm to solve

the unstructured search problem

i.e., find x_0 such that $f(x_0) = 1$

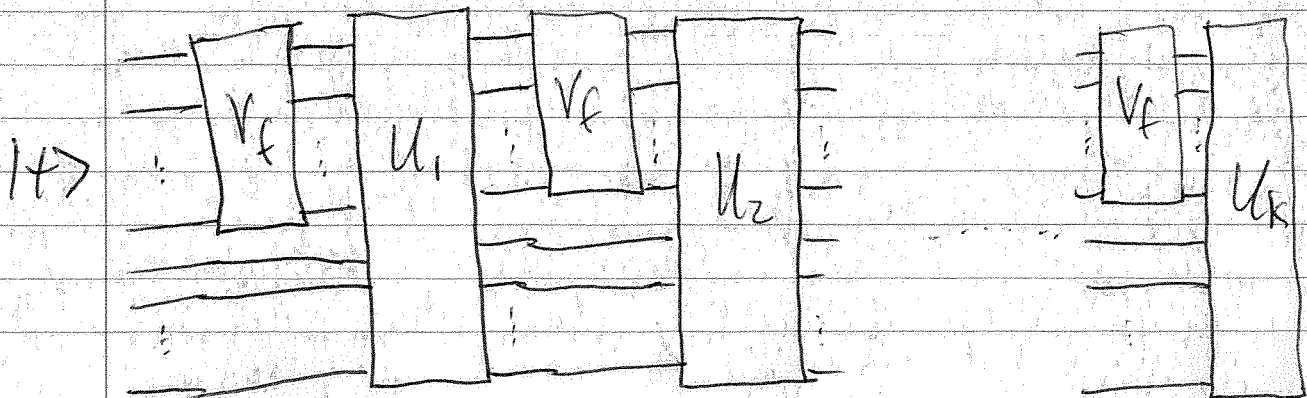
WLOG, let us assume that we
are using the phase oracle

$$V_f |x\rangle = (-1)^{f(x)} |x\rangle$$

$$\text{so that } V_f = (I - |x_0\rangle\langle x_0|) - |x_0\rangle\langle x_0|$$

$$= I - 2|x_0\rangle\langle x_0|$$

The most general quantum algorithm
for this task has the form

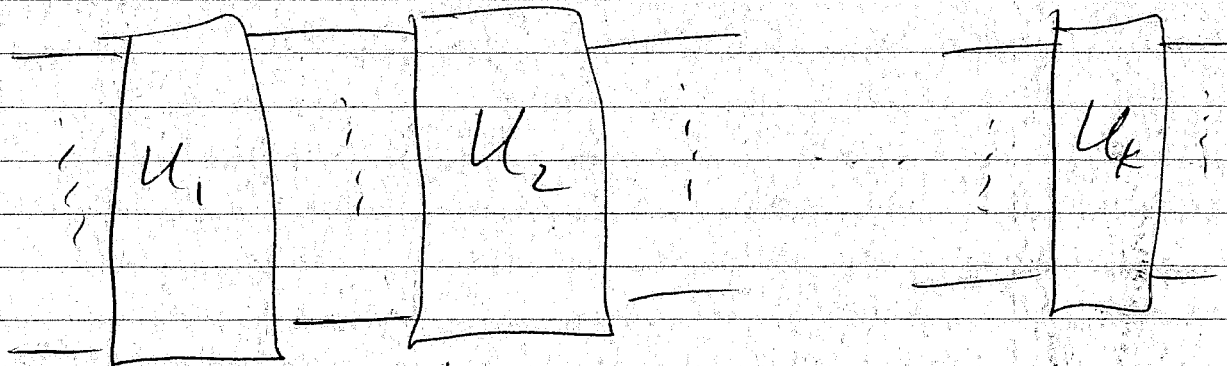


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Proof idea is ^{this} ~~to compare~~

If we do not call the oracle at all, then we don't learn any information about f & so we are no better off than essentially a random guess. Thus, if we compare the states after each U_k w/ & w/out the oracle calls, these states should become more & more orthogonal as the algorithm proceeds.

So we compare to



$$\text{Let } |\phi_{x_0}^k\rangle = \left(\prod_{i=1}^k U_i V_f \right) |\psi\rangle$$

$$|\phi^k\rangle = \left(\prod_{i=1}^k U_i \right) |\psi\rangle$$

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We could study an individual distance such as

$$\| |\phi_{x_0}^k\rangle - |\phi^k\rangle \|^2$$

but it turns out to be useful to symmetrize this expression by summing over all possible functions that have a single x_0 such that $f(x_0) = 1$. So we instead analyze

$$r_k = \sum_{x_0 \in \{0, 1\}^n} \| |\phi_{x_0}^k\rangle - |\phi^k\rangle \|^2$$

The idea will be

1) place an upper bound on r_k :

$$r_k \leq 4k^2$$

2) Supposing that we wish to identify x_0 w/ probability $> 1/2$, then this will imply that

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Part 1: bound $r_k \leq 4k^2$

$$\| |\phi_{x_0}^{k+1}\rangle - |\phi^{k+1}\rangle \|^2 = \| U_{k+1} V_f |\phi_{x_0}^k\rangle - U_{k+1} |\phi^k\rangle \|^2$$

$$= \| V_f |\phi_{x_0}^k\rangle - |\phi^k\rangle \|^2$$

$$= \| V_f (|\phi_{x_0}^k\rangle - |\phi^k\rangle) + (V_f - I) |\phi^k\rangle \|^2$$

$$= \| V_f (|\phi_{x_0}^k\rangle - |\phi^k\rangle) - 2|x_0\rangle \langle x_0 | \phi^k \rangle \|^2$$

$$= \| V_f (|\phi_{x_0}^k\rangle - |\phi^k\rangle) + (-2 \langle x_0 | \phi^k \rangle) |x_0\rangle \|^2$$

$$\leq \left(\| V_f (|\phi_{x_0}^k\rangle - |\phi^k\rangle) \| + 2 |\langle x_0 | \phi^k \rangle| \right)^2$$

$$= \left(\| |\phi_{x_0}^k\rangle - |\phi^k\rangle \| + 2 |\langle x_0 | \phi^k \rangle| \right)^2$$

$$= \| |\phi_{x_0}^k\rangle - |\phi^k\rangle \|^2 + 4 \| |\phi_{x_0}^k\rangle - |\phi^k\rangle \| \cdot |\langle x_0 | \phi^k \rangle|$$

$$+ 4 |\langle x_0 | \phi^k \rangle|^2$$

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Summing this over all x_0 , we get that

$$r_{k+1} \leq r_k + 4 \sum_{x_0 \in \{0,1\}^n} \left\| |\phi_{x_0}^k\rangle - |\phi^k\rangle \right\| \cdot |\langle x_0 | \phi^k \rangle| + 4$$

$$\uparrow \text{ b/c } 4 \sum_{x_0} |\langle x_0 | \phi^k \rangle|^2 = 4$$

But now we use Cauchy-Schwarz to upper bound

$$\begin{aligned} & \sum_{x_0 \in \{0,1\}^n} \left\| |\phi_{x_0}^k\rangle - |\phi^k\rangle \right\| |\langle x_0 | \phi^k \rangle| \\ & \leq \left(\sum_{x_0} \left\| |\phi_{x_0}^k\rangle - |\phi^k\rangle \right\|^2 \right)^{1/2} \left(\sum_{x_0} |\langle x_0 | \phi^k \rangle|^2 \right)^{1/2} \\ & = r_k^{1/2} \end{aligned}$$

So,

$$r_{k+1} \leq r_k + 4 + 4\sqrt{r_k}$$

Now

$$\begin{aligned} r_1 &= \sum_{x_0} \left\| |\phi_{x_0}^1\rangle - |\phi^1\rangle \right\|^2 \\ &= \sum_{x_0} \left\| |u, v_f | \psi \rangle - |u, \psi \rangle \right\|^2 \\ &= \sum_{x_0} \left\| (v_f - I) | \psi \rangle \right\|^2 \\ &= \sum_{x_0} \left\| -2 |x_0\rangle \langle x_0 | \psi \rangle \right\|^2 \\ &= 4 \sum_{x_0} |\langle x_0 | \psi \rangle|^2 = 4 \end{aligned}$$

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So by the induction, $r_k \leq 4k^2$, we get

$$\begin{aligned} r_{k+1} &\leq 4k^2 + 4 + 4\sqrt{4k^2} \\ &= 4(k^2 + 1 + 2k) \\ &= 4(k+1)^2 \end{aligned}$$

So we conclude that $r_k \leq 4k^2$

Conclusion is that the oracle queries can only increase distinguishability by so much.

Part 2: $r_k \geq c \cdot 2^n$

Suppose we want the algorithm to be successful w/ prob. $\geq 1/2$

So this means that $|\langle x_0 | \phi_{x_0}^k \rangle|^2 \geq 1/2$
we should have

But this implies that

$$\begin{aligned} &\| |\phi_{x_0}^k\rangle - e^{i\varphi} |x_0\rangle \|^2 \\ &= (\langle \phi_{x_0}^k | - e^{-i\varphi} \langle x_0 |) (|\phi_{x_0}^k\rangle - e^{i\varphi} |x_0\rangle) \\ &= 1 + 1 - e^{-i\varphi} \langle x_0 | \phi_{x_0}^k \rangle - e^{i\varphi} \langle \phi_{x_0}^k | x_0 \rangle \\ &\leq 2 - \sqrt{2} \end{aligned}$$

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Now come back to r_k :

$$\begin{aligned}
r_k &= \sum_{x_0 \in \{0, \pm 1\}^n} \left\| |\phi_{x_0}^k\rangle - |\phi^k\rangle \right\|^2 \\
&= \sum_{x_0 \in \{0, \pm 1\}^n} \left\| |\phi_{x_0}^k\rangle - |x_0\rangle + |x_0\rangle - |\phi^k\rangle \right\|^2 \\
&\geq \sum_{x_0 \in \{0, \pm 1\}^n} \left(\left\| |\phi_{x_0}^k\rangle - |x_0\rangle \right\| - \left\| |\phi^k\rangle - |x_0\rangle \right\| \right)^2 \\
&= \sum_{x_0} \left\| |\phi_{x_0}^k\rangle - |x_0\rangle \right\|^2 - \sum_{x_0} \left\| |\phi^k\rangle - |x_0\rangle \right\|^2 \\
&\quad - 2 \sum_{x_0} \left\| |\phi_{x_0}^k\rangle - |x_0\rangle \right\| \cdot \left\| |\phi^k\rangle - |x_0\rangle \right\|
\end{aligned}$$

Apply Cauchy-Schwarz to get

$$\sum_{x_0} \left\| |\phi_{x_0}^k\rangle - |x_0\rangle \right\| \cdot \left\| |\phi^k\rangle - |x_0\rangle \right\|$$

$$\leq \left[\sum_{x_0} \left\| |\phi_{x_0}^k\rangle - |x_0\rangle \right\|^2 \cdot \sum_{x_0} \left\| |\phi^k\rangle - |x_0\rangle \right\|^2 \right]^{1/2}$$

So conclude that

$$r_k \geq \left(\sum_{x_0} \left\| |\phi_{x_0}^k\rangle - |x_0\rangle \right\|^2 \right)^{1/2} - \left(\sum_{x_0} \left\| |\phi^k\rangle - |x_0\rangle \right\|^2 \right)^{1/2}$$

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From before, we require that

$$\| |\phi_{x_0}^k\rangle - |x_0\rangle \|^2 \leq 2 - \sqrt{2}$$

so this means that

$$\sum_{x_0} \| |\phi_{x_0}^k\rangle - |x_0\rangle \|^2 \leq 2^n (2 - \sqrt{2})$$

$$r_k \geq \left(- \left[(2 - \sqrt{2}) 2^n \right]^{1/2} + \left[\sum_{x_0} \| |\phi_k\rangle - |x_0\rangle \|^2 \right]^{1/2} \right)$$

so we bound from below

$$\begin{aligned} \sum_{x_0} \| |\phi_k\rangle - |x_0\rangle \|^2 &= \sum_{x_0} \left(\langle \phi_k | - \langle x_0 | \right) \left(|\phi_k\rangle - |x_0\rangle \right) \\ &= \sum_{x_0} \left(\langle \phi_k | \phi_k \rangle + \langle x_0 | x_0 \rangle - \langle \phi_k | x_0 \rangle - \langle x_0 | \phi_k \rangle \right) \\ &\geq \sum_{x_0} \left(2 - 2 \langle \phi_k | x_0 \rangle \right) \end{aligned}$$

But

$$\begin{aligned} \sum_{x_0} \langle \phi_k | x_0 \rangle &\leq \sqrt{\left(\sum_x 1 \right) \left(\sum_x \langle x | \phi_k \rangle^2 \right)} \\ &= \sqrt{2^n} \end{aligned}$$

so

$$\geq 2 \cdot 2^n - 2 \cdot \sqrt{2^n} = 2 \cdot (2^n - \sqrt{2^n})$$

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plug in to find that

$$\begin{aligned}
n_k &\geq \left(- \left[(2-\sqrt{2}) 2^n \right]^{1/2} + \left[2 \cdot (2^n - \sqrt{2}^n) \right]^{1/2} \right)^2 \\
&= \left(\left[(2-\sqrt{2}) 2^n \right]^{1/2} - \left[2 \cdot (2^n - \sqrt{2}^n) \right]^{1/2} \right)^2 \\
&\geq \left(\left[(2-\sqrt{2}) (2^n - \sqrt{2}^n) \right]^{1/2} - \left[2 \cdot (2^n - \sqrt{2}^n) \right]^{1/2} \right)^2 \\
&= \left[(2-\sqrt{2})^{1/2} - \sqrt{2} \right]^2 \cdot 2^n - \sqrt{2}^n \\
&\approx c \cdot 2^n \text{ for large } n
\end{aligned}$$

So combining

$$n_k \leq 4k^2$$

$$\text{if } n_k \geq c \cdot 2^n,$$

we see that $k \geq \Theta(\sqrt{2}^n)$

to have a reasonable success probability.

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Different topic: simulation of quantum systems
(simple ideas)

In QM, time evolution of a q. state is governed by the Schrödinger equation:

$$i\hbar \frac{d}{dt} |\psi(t)\rangle = H(t) |\psi(t)\rangle$$

$H(t)$ is the Hamiltonian,
generator of time translations

usually take $\hbar=1$ for convenience

Also take $H(t)$ to be
time independent for ~~convenience~~ ^{simplicity}
so that

$$H(t) = H$$

\Rightarrow solution is $|\psi(t)\rangle = e^{-iHt} |\psi(0)\rangle$

Say that a Hamiltonian H
acting on n qubits can be efficiently
simulated if $\forall t, \epsilon > 0$

\exists a quantum circuit U consisting of
 $\text{poly}(n, t, 1/\epsilon)$ gates such that $\|U - e^{-iHt}\| \leq \epsilon$

we will not be able to simulate arbitrary Hamiltonians, but some of them we can.

Some trivial cases: if H acts on only a constant # of qubits then we just invoke Solovay-Kitaev to say there exists a quantum circuit of error at most ϵ using $\text{poly}(\log(\frac{1}{\epsilon}))$ 1- & 2-qubit gates

can also rotate the basis in which H is applied using any U w/ an efficient decomposition into basic gates.

That is, if H has an efficient sim. & U has an eff. implementation, then UHU^\dagger can be efficiently simulated. Follows from

$$Ue^{-iHt}U^\dagger = e^{-iUHU^\dagger t}$$

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Suppose a Hamiltonian is diagonal
in the computational basis $|a\rangle$
that any diagonal element can
be computed efficiently

$$d(a) = \langle a | H | a \rangle$$

then we just do

$$|a, 0\rangle \rightarrow |a, d(a)\rangle \quad \text{compute}$$

$$\rightarrow e^{-itd(a)} |a, d(a)\rangle \quad \text{add phase}$$

$$\rightarrow e^{-itd(a)} |a, 0\rangle \quad \text{uncompute}$$

$$= e^{-itH} |a, 0\rangle$$

extend by linearity

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Many physical Hamiltonians are expressed as a sum of terms

For example, particle in a potential

$$H = \frac{p^2}{2m} + V(x)$$

idea is to discretize the x coordinate

$V(x)$ is diagonal of natural discretizations of $p^2 = -\frac{d^2}{dx^2}$

are diagonal in the discrete Fourier basis.

so we can efficiently simulate both $V(x)$ + $p^2/2m$

Another example:

Ham. for spin system

$$H = \sum_i h_i X_i + \sum_{\langle ij \rangle} J_{ij} Z_i Z_j$$

This is a sum of terms, each of which acts on no more than 2 qubits + thus are each easy to simulate

If H_1 & H_2 can be efficiently simulated, then so can

$$H_1 + H_2$$

How is this so?

If they commute, obvious

$$\text{b/c } e^{-iH_1 t} e^{-iH_2 t} = e^{-i(H_1 + H_2)t}$$

But if they don't?

We have the Lie-Trotter product formula:

$$e^{-i(H_1 + H_2)t} = \lim_{m \rightarrow \infty} \left(e^{-\frac{iH_1 t}{m}} e^{-\frac{iH_2 t}{m}} \right)^m$$

proof idea:

Taylor expand the exponential

can truncate the RHS to a finite # of terms & if we want to have

$$\left\| \left(e^{-\frac{iH_1 t}{m}} e^{-\frac{iH_2 t}{m}} \right)^m - e^{-i(H_1 + H_2)t} \right\|$$

we need $m = O((vt)^2 / \epsilon) \leq \epsilon$
where $v = \max \{ \|H_1\|, \|H_2\| \}$

can get improvements (smaller m)
by considering

$$\left(e^{-\frac{iH_1 t}{2m}} e^{-\frac{iH_2 t}{m}} e^{-\frac{iH_1 t}{2m}} \right)^m$$

In general, we have

$$e^{-i(H_1 + \dots + H_k)t} =$$

$$\lim_{m \rightarrow \infty} \left(e^{-\frac{iH_1 t}{m}} \dots e^{-\frac{iH_k t}{m}} \right)^m$$