

Lecture 18

①

An important state in Gaussian QI is the two-mode squeezed vacuum state.

It is generated by the action of a two-mode squeezing transformation on a two-mode vacuum state:

$$\hat{S}_r |0\rangle|0\rangle$$

where

$$\hat{S}_r = e^{-\frac{i}{2} \hat{r}^T (r H_{TMS}) \hat{r}}$$

where $H_{TMS} = \begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix}$ & $r > 0$ is squeezing strength

so that

$$\hat{H}_{TMS} = i r (\hat{a}_1^\dagger \hat{a}_2^\dagger - \hat{a}_1 \hat{a}_2) \quad \& \quad \hat{S}_r = e^{-i \hat{H}_{TMS}}$$

& symplectic transformation is

(3)

Given that $\cosh r \geq 1 \quad \forall r \in \mathbb{R}$

& cov. matrix B of I_2

this is the covariance matrix for
a thermal state.

We can set $\cosh(2r) = 2\bar{n} + 1$

& this implies that $\sinh r = 2\sqrt{\bar{n}(\bar{n}+1)}$

\Rightarrow cov. matrix for TMSV is

$$\begin{bmatrix} (2\bar{n}+1)I_2 & 2\sqrt{\bar{n}(\bar{n}+1)}\sigma_z \\ 2\sqrt{\bar{n}(\bar{n}+1)}\sigma_z & (2\bar{n}+1)I_2 \end{bmatrix}$$

In Hilbert space, the state

is given by

$$|\Psi_r\rangle = e^{-i\hat{H}_{\text{TMS}}} |0\rangle|0\rangle = e^{r(\hat{a}_1^\dagger \hat{a}_2^\dagger - \hat{a}_1 \hat{a}_2)} |0\rangle|0\rangle$$

(7)

can find this by means of
disentangling theorem, itself a
consequence of BCH:

$$e^{r(\hat{a}_1^\dagger \hat{a}_2^\dagger - \hat{a}_1 \hat{a}_2)}$$
$$= e^{\Gamma \hat{a}_1^\dagger \hat{a}_2^\dagger} e^{-g(\hat{a}_1^\dagger \hat{a}_1 + \hat{a}_2^\dagger \hat{a}_2 + 1)}$$
$$e^{-\Gamma \hat{a}_1 \hat{a}_2}$$

where $\Gamma = \tanh(r)$

$$g = \ln(\cosh(r))$$

Consider that

$$e^{-\Gamma \hat{a}_1 \hat{a}_2} |0\rangle |0\rangle$$
$$= \sum_{l=0}^{\infty} \frac{(-\Gamma \hat{a}_1 \hat{a}_2)^l}{l!} |0\rangle |0\rangle = |0\rangle |0\rangle$$

then $e^{-g(\hat{a}_1^\dagger \hat{a}_1 + \hat{a}_2^\dagger \hat{a}_2 + 1)} |0\rangle |0\rangle$

$$= e^{-g(\hat{a}_1^\dagger \hat{a}_1)} e^{-g(\hat{a}_2^\dagger \hat{a}_2)} e^{-g} |0\rangle |0\rangle$$

(5)

$$\begin{aligned} &= e^{-g} |0\rangle|0\rangle \quad (\text{since } e^{-g \hat{a}_1^\dagger \hat{a}_2} |0\rangle \\ &= \frac{1}{\cosh r} |0\rangle|0\rangle \quad = e^{-g \cdot 0} |0\rangle \\ & \quad \quad \quad \quad \quad \quad \quad = |0\rangle) \end{aligned}$$

Finally,

$$\begin{aligned} &e^{\Gamma \hat{a}_1^\dagger \hat{a}_2} |0\rangle|0\rangle \\ &= \sum_{l=0}^{\infty} \frac{(\Gamma \hat{a}_1^\dagger \hat{a}_2)^l}{l!} |0\rangle|0\rangle \\ &= \sum_{l=0}^{\infty} (\tanh r)^l \frac{\hat{a}_1^\dagger}{\sqrt{l!}} \frac{\hat{a}_2}{\sqrt{l!}} |0\rangle|0\rangle \\ &= \sum_{l=0}^{\infty} (\tanh r)^l |l\rangle|l\rangle \end{aligned}$$

$$\Rightarrow |\Psi_r\rangle = \frac{1}{\cosh r} \sum_{n=0}^{\infty} (\tanh r)^n |n\rangle|n\rangle$$

(6)

For the other parametrisation,
we have

$$\frac{1}{\cosh r} = \frac{1}{\sqrt{\bar{n}+1}}$$

$$\& \tanh r = \sqrt{\frac{\bar{n}}{\bar{n}+1}}$$

$$\Rightarrow |4_r\rangle = \frac{1}{\sqrt{\bar{n}+1}} \sum_{n=0}^{\infty} \left(\sqrt{\frac{\bar{n}}{\bar{n}+1}} \right)^n |n\rangle |n\rangle$$

Gaussian quantum channels

previously, we discussed Gaussian unitaries. Now we consider channels that map Gaussian states to Gaussian states.

- Gaussian channels maps an n -mode state to an m -mode state is characterized by

(7)

by $2m \times 2m$ ^{real} matrix X ,
called scaling matrix, \dagger

$2m \times 2m$ real symmetric matrix
 Y , called noise matrix,
 \dagger a displacement S where
 $S \in \mathbb{R}^{2m}$.

The effect of the channel on
an input Gaussian state w/
mean vector $\bar{r} \in \mathbb{R}^{2m}$ \dagger cov.
matrix $\sigma \in \mathbb{R}^{2m \times 2m}$ is

$$\bar{r} \rightarrow X\bar{r} + S$$

$$\sigma \rightarrow X\sigma X^T + Y$$

X \dagger Y should satisfy

$$Y + i\Omega_m \geq iX\Omega_n X^T$$

8

This condition comes about
in order for the output state
of the channel to be a legitimate
state. (respecting the uncertainty
relation $\sigma + i\Omega \geq 0$)

We can prove one direction:

If X & Y correspond to a Gaussian
channel, then $Y + i\Omega \geq iX\Omega X^T$.

Consider the action of the channel

$$N \text{ on } |\psi_{r,n}\rangle_{RA} = |\psi_r\rangle^{\otimes n}$$

$$(\text{id}_R \otimes N_{A \rightarrow B}) (|\psi_{r,n}\rangle \langle \psi_{r,n}|_{RA})$$

(A is n-mode system)

Consider that cov. matrix of $|\psi_{r,n}\rangle_{RA}$ is

$$\sigma_n(r) = \begin{bmatrix} \cosh(2r) I_{2n} & \sinh(2r) \Sigma_n \\ \sinh(2r) \Sigma_n & \cosh(2r) I_{2n} \end{bmatrix}$$

where $\Sigma_n = \bigoplus_{j=1}^n \sigma_z = I_{2n} \otimes \sigma_z$ (9)

Then covariance matrix of output state is

$$\begin{aligned} & \begin{bmatrix} I_{2n} & 0 \\ 0 & X \end{bmatrix} \begin{bmatrix} \sigma_n(r) \end{bmatrix} \begin{bmatrix} I_{2n} & 0 \\ 0 & X \end{bmatrix} + \begin{bmatrix} 0_{2n} & 0 \\ 0 & Y \end{bmatrix} \\ &= \begin{bmatrix} \cosh(2r) I_{2n} & \sinh(2r) \Sigma_n X^T \\ \sinh(2r) X \Sigma_n & \cosh(2r) X X^T + Y \end{bmatrix} \\ &= \sigma_{\text{out}}(r) \end{aligned}$$

Given that N is a channel,

this is a legitimate covariance matrix, satisfying

$$\sigma_{\text{out}}(r) + i\mathcal{L}_{n+m} \geq 0$$

or

$$\begin{bmatrix} \cosh(2r) I_{2n} + i\mathcal{L}_n & \sinh(2r) \Sigma_n X^T \\ \sinh(2r) X \Sigma_n & \cosh(2r) X X^T + Y + i\mathcal{L}_m \end{bmatrix} \geq 0$$

(10)

Now invoke Schur complement lemma:

$$\mathbb{Z} = \begin{bmatrix} A & B \\ B^T & C \end{bmatrix} > 0 \Leftrightarrow$$

$$A > 0 \quad \& \quad C - B^T A^{-1} B > 0$$

(holds for pos. def., but can be generalized to PSD case)

So previous condition is equivalent to

$$\cosh(2r) I_{2n} + i \Omega_n > 0 \quad \&$$

$$(*) \quad \cosh(2r) X X^T + Y + i \Omega_m$$

$$- \sinh(2r) X \Sigma_n \left(\cosh(2r) \frac{I_{2n} + i \Omega_n}{\cosh(2r)} \right)^{-1}$$

$$\sinh(2r) \Sigma_n X^T > 0$$

(11)

Consider that

$$\begin{bmatrix} \cosh(2r)I_2 \\ + i\Omega_1 \end{bmatrix}^{-1} = \begin{bmatrix} \cosh(2r) & i \\ -i & \cosh(2r) \end{bmatrix}^{-1}$$

$$= \frac{\cosh(2r)I_2 - i\Omega_1}{\sinh^2(2r)}$$

$$\Rightarrow (*) = \cosh(2r)XX^T + Y + i\Omega_m$$

$$- \sinh^2(2r) X \sum_n \left(\frac{\cosh(2r)I_{2n} - i\Omega_n}{\sinh^2(2r)} \right) \sum_n X^T$$

$$= \cosh(2r)XX^T + Y + i\Omega_m > 0$$

$$(**) - X \sum_n (\cosh(2r)I_{2n} - i\Omega_n) \sum_n X^T > 0$$

~~Consider that 2nd term~~

$$= ~~0~~ - \cosh(2r)XX^T + iX \sum_n \Omega_n \sum_n X^T$$

~~0~~

(12)

$$\Rightarrow = -\cosh(2r) X X^T$$

$$\begin{aligned} & + i X (I_n \otimes \sigma_z) (I_n \otimes \Lambda_1) (I_n \otimes \sigma_z) X^T \\ & = -\cosh(2r) X X^T \quad (\text{using } \sigma_z \Lambda_1 \sigma_z = -\Lambda_1) \\ & - i X \Lambda_n X^T \end{aligned}$$

$$\Rightarrow (***) = Y + i \Lambda_m$$

$$- i X \Lambda_n X^T > 0$$

$$\Rightarrow Y + i \Lambda_m > i X \Lambda_n X^T$$

Then invoking Schur comp. generalization for PSD, we get

$$Y + i \Lambda_m \geq i X \Lambda_n X^T$$

(13)

Wigner function of a quantum channel

$$W_{\mathcal{N}}(r'|r) = (2\pi)^n \text{Tr} \left[\hat{A}_{r'} \mathcal{N}(\hat{A}_r) \right]$$

where $\hat{A}_r = \hat{D}_r \hat{A}_0 \hat{D}_r$

$$\hat{A}_0 = \frac{1}{(2\pi)^n} \int dr' \hat{D}_{-r'}$$

Using definitions

this then reduces to

$$\frac{(2\pi)^n}{(2\pi)^{4n}} \int dr'' \int dr''' e^{i(r'')^T \mathcal{N} r} e^{i(r''')^T \mathcal{N} r'} \text{Tr} \left[\hat{D}_{-r'''} \mathcal{N}(\hat{D}_{-r''}) \right]$$

Using $\text{Tr} \left[\hat{D}_{-r'''} \mathcal{N}(\hat{D}_{-r''}) \right]$

$$= \text{Tr} \left[\mathcal{N}^\dagger(\hat{D}_{-r'''}) \hat{D}_{-r''} \right]$$

$$\downarrow \mathcal{N}^\dagger(\hat{D}_z) = \hat{D}_z \mathcal{N}^\dagger \mathcal{N}^\dagger z e^{-\frac{1}{4} z^T \mathcal{N} \mathcal{N}^\dagger z + i z^T \mathcal{N} z}$$

(14)

for Gaussian channel X characterized
by scaling Λ , noise Y , & shift δ

$$\Rightarrow W_X(r'|r) =$$

$$\frac{(2\pi)^m}{(2\pi)^{4n}} \iint dr'' dr''' e^{i[(r'')^T \Lambda r + (r''')^T \Lambda r']}$$

$$\text{Tr} \left[\hat{D}_{-\Lambda X^T \Lambda^T r'''} \hat{D}_{-r''} \right]$$

$$e^{-\frac{1}{4} (r''')^T \Lambda Y \Lambda^T r'''} - i (r''')^T \Lambda \delta$$

$$\text{using } \text{Tr} \left[\hat{D}_{r_1} \hat{D}_{-r_2} \right] \\ = (2\pi)^2 \delta^{(2n)}(r_1 - r_2)$$

$$\Rightarrow W_X(r'|r) =$$

$$\frac{1}{(2\pi)^{2n}} \iint dr'' dr''' e^{i[(r'')^T \Lambda r + (r''')^T \Lambda r']} \\ \delta^{(2n)}(r'' + \Lambda X^T \Lambda^T r''')$$

$$e^{-\frac{1}{4} (r''')^T \Lambda Y \Lambda^T r'''} - i (r''')^T \Lambda \delta$$

(15)

$$= \frac{1}{(2\pi)^{2n}} \int dr''' e^{i[-(\Omega X^T \Omega^T r''')^T \Omega r - (r')^T \Omega r''']} e^{-\frac{1}{4}(r''')^T \Omega Y \Omega^T r'''} + i \delta^T \Omega r'''$$

$$= \frac{1}{(2\pi)^{2n}} \int dr''' e^{-i[(r''')^T \Omega X r - (r')^T \Omega r''']} e^{-\frac{1}{4}(r''')^T \Omega Y \Omega^T r'''} + i \delta^T \Omega r'''$$

$$= \frac{1}{(2\pi)^{2n}} \int dr''' e^{i([Xr - r' + \delta]^T \Omega r''')} e^{-\frac{1}{4}(r''')^T \Omega Y \Omega^T r'''} \text{rewrite again as}$$

$$= \frac{1}{(2\pi)^{2n}} \int dr''' e^{(r''')^T [\Omega^T i(Xr - r' + \delta)]} e^{-(r''')^T (\frac{1}{4} \Omega Y \Omega^T) r'''} \text{rewrite again as}$$

$$A = \frac{1}{4} \Omega Y \Omega^T$$

$$b = i \Omega^T (Xr - r' + \delta)$$

(16)

Use Gaussian integration
formula

$$= \frac{1}{(2\pi)^{2n}} \frac{\pi^n}{\sqrt{\text{Det}(\frac{1}{4} \mathcal{N} Y \mathcal{N}^T)}} \times$$

$$e^{\frac{1}{4} (i \mathcal{N}^T (x_r - r' + s))^T (\frac{1}{4} \mathcal{N} Y \mathcal{N}^T)^{-1} (i \mathcal{N}^T (x_r - r' + s))}$$

$$= \frac{1}{\pi^n \sqrt{\text{Det}(Y)}} e^{- (x_r - r' + s)^T Y^{-1} (x_r - r' + s)}$$

then conclude that

$$W_{Xr}(r'|r) =$$

$$\frac{1}{\pi^n \sqrt{\text{Det}(Y)}} e^{- (r' - (x_r + s))^T Y^{-1} (r' - (x_r + s))}$$

this is the same as the
conditional prob. dist. for a
classical Gaussian channel