

Lecture 13

①

Characteristic functions of
quasi-probability distributions

recall that for a single mode

$$\hat{D}_{-r} \equiv e^{-ir^\dagger \hat{a} r}$$

with convention

$$\hat{D}_\alpha \equiv e^{\alpha \hat{a}^\dagger - \alpha^* \hat{a}}$$

$$\text{with } \alpha = \frac{x + ip}{\sqrt{2}} \quad \text{with } r = \begin{bmatrix} x \\ p \end{bmatrix}$$

then

$$\hat{D}_{-r} = \hat{D}_\alpha$$

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Then define coherent states as

$$\hat{D}_\alpha |0\rangle = |\alpha\rangle = e^{-\frac{|\alpha|^2}{2}} \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} |n\rangle$$

recall that

$$\hat{D}_\alpha \hat{D}_\beta = e^{\frac{1}{2}(\alpha\beta^* - \alpha^*\beta)} \hat{D}_{\alpha+\beta}$$

$$\begin{aligned} \Rightarrow \langle \beta | \alpha \rangle &= \langle 0 | \hat{D}_{-\beta} \hat{D}_\alpha | 0 \rangle \\ &= \langle 0 | \hat{D}_{\alpha-\beta} | 0 \rangle e^{\frac{1}{2}(\alpha\beta^* - \alpha^*\beta)} \\ &= \langle 0 | \alpha - \beta \rangle e^{\frac{1}{2}(\alpha\beta^* - \alpha^*\beta)} \\ &= e^{-\frac{1}{2}|\alpha-\beta|^2} e^{\frac{1}{2}(\alpha\beta^* - \alpha^*\beta)} \end{aligned}$$

Coherent states form an overcomplete set:

$$\frac{1}{\pi} \int_{\mathbb{C}} d^2\alpha |\alpha\rangle \langle \alpha| = \mathbb{I}$$

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can then use coherent-state basis
to evaluate traces of trace-class
operators:

$$\begin{aligned}\text{Tr}[\hat{O}] &= \sum_{m=0}^{\infty} \langle m | \hat{O} | m \rangle \\ &= \sum_{m=0}^{\infty} \langle m | \frac{1}{\pi} \int_{\mathbb{C}} d^2\alpha |\alpha\rangle \langle \alpha | \hat{O} | m \rangle \\ &= \frac{1}{\pi} \int_{\mathbb{C}} d^2\alpha \sum_{m=0}^{\infty} \langle m | \alpha \rangle \langle \alpha | \hat{O} | m \rangle \\ &= \frac{1}{\pi} \int_{\mathbb{C}} d^2\alpha \sum_{m=0}^{\infty} \langle \alpha | \hat{O} | m \rangle \langle m | \alpha \rangle \\ &= \frac{1}{\pi} \int_{\mathbb{C}} d^2\alpha \langle \alpha | \hat{O} | \alpha \rangle\end{aligned}$$

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can use this to easily figure out
mean vector & covariance matrix.

Alternatively, use that (even simpler)

$$\begin{aligned} & \text{Tr}[\hat{x} |\alpha\rangle\langle\alpha|] \\ &= \langle\alpha|\hat{x}|\alpha\rangle \\ &= \langle\alpha|\frac{\hat{a} + \hat{a}^\dagger}{\sqrt{2}}|\alpha\rangle = \frac{\langle\alpha|\hat{a}|\alpha\rangle + \langle\alpha|\hat{a}^\dagger|\alpha\rangle}{\sqrt{2}} \\ &= \frac{\langle\alpha|\alpha|\alpha\rangle + \langle\alpha|\alpha^*|\alpha\rangle}{\sqrt{2}} \\ &= \frac{2 \text{Re}\{\alpha\}}{\sqrt{2}} = \sqrt{2} \text{Re}\{\alpha\} \end{aligned}$$

Similarly,

$$\langle\alpha|\hat{p}|\alpha\rangle = \sqrt{2} \text{Im}\{\alpha\}$$

Alternatively,

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$$\begin{aligned}\text{Tr}[\hat{x}|\alpha\rangle\langle\alpha|] &= \\ \text{Tr}[\hat{x} \hat{D}_\alpha |0\rangle\langle 0| \hat{D}_{-\alpha}] &= \\ = \text{Tr}[\hat{D}_{-\alpha} \hat{x} \hat{D}_\alpha |0\rangle\langle 0|] &= \\ = \text{Tr}[\hat{D}_r \hat{x} \hat{D}_{-r} |0\rangle\langle 0|] &= \\ = \text{Tr}[(\hat{x} + x) |0\rangle\langle 0|] &= \\ = \langle 0|\hat{x}|0\rangle + x = x = \sqrt{2} \text{Re}\{\alpha\} &\end{aligned}$$

For covariance matrix,

$$\begin{aligned}\text{Tr}[\hat{x}^2 |\alpha\rangle\langle\alpha|] &= \\ = \text{Tr}[\hat{x}^2 \hat{D}_\alpha |0\rangle\langle 0| \hat{D}_{-\alpha}] &= \\ = \text{Tr}[\hat{D}_r \hat{x}^2 \hat{D}_{-r} |0\rangle\langle 0|] &= \\ = \text{Tr}[(\hat{x} + x)^2 |0\rangle\langle 0|] &= \\ = \text{Tr}[(\hat{x}^2 + 2x\hat{x} + x^2) |0\rangle\langle 0|] &\end{aligned}$$

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$$= 1 + x^2$$

$$\Rightarrow \text{Tr} [(\hat{x} - x)^2 |\alpha\rangle\langle\alpha|]$$

$$= \underline{\underline{1}}$$

keep calculating ↓ cov. matrix of

$$|\alpha\rangle\langle\alpha| \text{ is } \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

(same as vacuum)

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Trace of a displacement operator
in terms of Dirac delta function
(main idea upon which
characteristic functions are
built)

Consider

$$\langle \alpha | D(\beta) | \alpha \rangle$$

$$= \langle 0 | D^\dagger(\alpha) D(\beta) D(\alpha) | 0 \rangle$$

$$= \langle 0 | D^\dagger(\alpha) \exp[\beta \hat{a}^\dagger - \beta^* \hat{a}] D(\alpha) | 0 \rangle$$

$$= \langle 0 | \exp[D^\dagger(\alpha) (\beta \hat{a}^\dagger - \beta^* \hat{a}) D(\alpha)] | 0 \rangle$$

$$= \langle 0 | \exp[\beta (\hat{a}^\dagger + \alpha) - \beta^* (\hat{a} + \alpha)] | 0 \rangle$$

$$= e^{\beta \alpha^* - \beta^* \alpha} \langle 0 | \exp[\beta \hat{a}^\dagger - \beta^* \hat{a}] | 0 \rangle$$

$$= e^{\beta \alpha^* - \beta^* \alpha} \langle 0 | D(\beta) | 0 \rangle$$

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$$= e^{\beta \alpha^* - \beta^* \alpha} \langle 0 | \beta \rangle$$

$$= e^{\beta \alpha^* - \beta^* \alpha} e^{-\frac{1}{2} |\beta|^2}$$

$$\text{Then } \text{Tr}[D(\beta)] = \frac{1}{\pi} \int_{\mathbb{C}} d^2 \alpha \langle \alpha | D(\beta) | \alpha \rangle$$

$$= \frac{1}{\pi} e^{-\frac{1}{2} |\beta|^2} \int_{\mathbb{C}} d^2 \alpha e^{\beta \alpha^* - \beta^* \alpha}$$

$$\text{set } \alpha = x + iy$$

$$\beta = u + iv$$

$$\begin{aligned} \Rightarrow \beta \alpha^* - \beta^* \alpha &= (u + iv)(x - iy) \\ &\quad - (u - iv)(x + iy) \\ &= i[2vx - 2uy] \end{aligned}$$

$$\begin{aligned} \Rightarrow \int_{\mathbb{C}} d^2 \alpha e^{\beta \alpha^* - \beta^* \alpha} &= \iint dx dy e^{i[2vx - 2uy]} \\ &= \int dx e^{i2vx} \cdot \int dy e^{-i2uy} \end{aligned}$$

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$$= 2\pi \delta(2v) \cdot 2\pi \delta(2u)$$

$$= \pi \delta(v) \cdot \pi \delta(u)$$

(used $\delta(ax) = \frac{\delta(x)}{|a|}$)

$$\Rightarrow \text{Tr} [D(\beta)] = \frac{1}{\pi} e^{-\frac{1}{2}|\beta|^2} \cdot \pi^2 \delta^2(\beta)$$

$$= \pi \delta^2(\beta)$$

since $e^{-\frac{1}{2}|\beta|^2} \Big|_{\beta=0} = 1$

then by linearity

$$\text{Tr} [\hat{D}_\alpha \hat{D}_\beta] = \text{Tr} [\hat{D}_{\alpha-\beta}] \times$$

$$e^{\frac{1}{2}(-\alpha\beta^* + \alpha^*\beta)}$$

$$= \pi \delta^2(\alpha-\beta) e^{\frac{1}{2}(-\alpha\beta^* + \alpha^*\beta)}$$

$$= \pi \delta^2(\alpha-\beta)$$

since $\delta^2(\alpha-\beta) = 0$ when $\alpha \neq \beta$

$$\downarrow e^{\frac{1}{2}(-\alpha\beta^* + \alpha^*\beta)} = 1 \text{ when } \alpha = \beta$$

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⇒ trace orthogonality relationship is

$$\text{Tr}[\hat{D}_\alpha \hat{D}_{-\beta}] = \pi \delta^2(\alpha - \beta)$$

then for $\alpha = \frac{x_r + ip_r}{\sqrt{2}}$ $\beta = \frac{x_s + ip_s}{\sqrt{2}}$

$$\text{Tr}[\hat{D}_r \hat{D}_{-s}] =$$

$$\pi \delta^2(\alpha - \beta)$$

$$= \pi \delta\left(\frac{x_r - x_s}{\sqrt{2}}\right) \cdot \delta\left(\frac{p_r - p_s}{\sqrt{2}}\right)$$

$$= 2\pi \delta(x_r - x_s) \cdot \delta(p_r - p_s)$$

$$= 2\pi \delta^2(r - s)$$

$$\text{So } \text{Tr}[\hat{D}_r \hat{D}_{-s}] = 2\pi \delta^2(r - s)$$

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for n modes

$$\text{Tr} [\hat{D}_{\underline{\alpha}} \hat{D}_{-\underline{\beta}}] = \pi^n \delta^{2n}(\underline{\alpha} - \underline{\beta})$$

$$\text{Tr} [\hat{D}_{\underline{r}} \hat{D}_{-\underline{s}}] = (2\pi)^n \delta^{2n}(\underline{r} - \underline{s})$$

Characteristic function for a

trace-class operator ρ is defined as

$$\chi_{\rho}(\alpha) = \text{Tr} [\hat{D}_{\alpha} \rho]$$

this is finite $\forall \alpha \in \mathbb{C}^{2n}$ b/c

$$\begin{aligned} |\text{Tr} [\hat{D}_{\alpha} \rho]| &\leq \|\hat{D}_{\alpha}\|_{\infty} \|\rho\|_1 \\ &= \|\rho\|_1 < \infty \end{aligned}$$

Hölder inequality.

can then write

$$\rho = \int d^2\alpha \chi_\rho(\alpha) \hat{D}_{-\alpha}$$

~~Consider that~~

Proof:

$$\text{Tr}[\hat{D}_\beta \rho] =$$

$$\begin{aligned} \text{Since } \rho &= \left[\frac{1}{\pi} \int d^2\alpha |\alpha\rangle\langle\alpha| \right] \rho \left[\frac{1}{\pi} \int d^2\beta |\beta\rangle\langle\beta| \right] \\ &= \frac{1}{\pi^2} \iint d^2\alpha \langle\alpha|\rho|\beta\rangle |\alpha\rangle\langle\beta| \end{aligned}$$

it suffices to show that

$$|\alpha\rangle\langle\beta| = \frac{1}{\pi} \int d^2\gamma \text{Tr} [|\alpha\rangle\langle\beta| \hat{D}_\gamma] \hat{D}_{-\gamma}$$

This is the same as

$$\begin{aligned} |0\rangle\langle 0| &= \frac{1}{\pi} \int d^2\gamma \text{Tr} [|\alpha\rangle\langle\beta| \hat{D}_\gamma] \hat{D}_{-\alpha} \hat{D}_{-\gamma} \hat{D}_\beta \\ &= \frac{1}{\pi} \int d^2\gamma \langle\beta-\gamma|\alpha\rangle e^{\frac{1}{2}(\gamma\beta^* - \gamma^*\beta)} \hat{D}_{-\alpha} \hat{D}_{-\gamma} \hat{D}_\beta \end{aligned}$$

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$$= \frac{1}{\pi} \int d^2y \ e^{-\frac{1}{2}|\beta-\alpha-y|^2} \hat{D}_{\beta-\alpha-y}$$

$$= \frac{1}{\pi} \int d^2y \ e^{-\frac{1}{2}|y|^2} \hat{D}_y$$

pick up next time.