

Recoverability in quantum information theory

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Main message

- Entropy inequalities established in the 1970s are a mathematical consequence of the postulates of quantum physics
- They are helpful in determining the ultimate limits on many physical processes (communication, thermodynamics, uncertainty relations)
- Many of these entropy inequalities are equivalent to each other, so we can say that together they constitute a fundamental law of quantum information theory
- There has been recent interest in refining these inequalities, trying to understand how well one can attempt to reverse an irreversible physical process
- This poster presentation discusses progress in this direction

Umegaki relative entropy [Ume62]

The quantum relative entropy is a measure of dissimilarity between two quantum states. Defined for state ρ and positive semi-definite σ as

$$D(\rho\|\sigma) \equiv \text{Tr}\{\rho[\log \rho - \log \sigma]\}$$

whenever $\text{supp}(\rho) \subseteq \text{supp}(\sigma)$ and $+\infty$ otherwise

Physical interpretation with quantum Stein's lemma [HP91, NO00]

Given are n quantum systems, all of which are prepared in either the state ρ or σ . With a constraint of $\varepsilon \in (0, 1)$ on the Type I error of misidentifying ρ , then the optimal error exponent for the Type II error of misidentifying σ is $D(\rho\|\sigma)$.

Relative entropy as “mother” entropy

Many important entropies can be written in terms of relative entropy:

- $H(A)_\rho \equiv -D(\rho_A \| I_A)$ (entropy)
- $H(A|B)_\rho \equiv -D(\rho_{AB} \| I_A \otimes \rho_B)$ (conditional entropy)
- $I(A; B)_\rho \equiv D(\rho_{AB} \| \rho_A \otimes \rho_B)$ (mutual information)
- $I(A; B|C)_\rho \equiv D(\rho_{ABC} \| \exp\{\log \rho_{AC} + \log \rho_{BC} - \log \rho_C\})$ (cond. MI)

Equivalences

- $H(A|B)_\rho = H(AB)_\rho - H(B)_\rho$
- $I(A; B)_\rho = H(A)_\rho + H(B)_\rho - H(AB)_\rho$
- $I(A; B)_\rho = H(B)_\rho - H(B|A)_\rho$
- $I(A; B|C)_\rho = H(AC)_\rho + H(BC)_\rho - H(ABC)_\rho - H(C)_\rho$
- $I(A; B|C)_\rho = H(B|C)_\rho - H(B|AC)_\rho$

Monotonicity of quantum relative entropy [Lin75, Uhl77]

Let ρ be a state, let σ be positive semi-definite, and let \mathcal{N} be a quantum channel. Then

$$D(\rho\|\sigma) \geq D(\mathcal{N}(\rho)\|\mathcal{N}(\sigma))$$

“Distinguishability does not increase under a physical process”

Characterizes a fundamental irreversibility in any physical process

Proof approaches

- Lieb concavity theorem [L73]
- relative modular operator method (see, e.g., [NP04])
- quantum Stein’s lemma [BS03]

Strong subadditivity

Strong subadditivity [LR73]

Let ρ_{ABC} be a tripartite quantum state. Then

$$I(A; B|C)_\rho \geq 0$$

Equivalent statements (by definition)

- Entropy sum of two individual systems is larger than entropy sum of their union and intersection:

$$H(AC)_\rho + H(BC)_\rho \geq H(ABC)_\rho + H(C)_\rho$$

- Conditional entropy does not decrease under the loss of system A :

$$H(B|C)_\rho \geq H(B|AC)_\rho$$

When does equality in monotonicity of relative entropy hold?

- $D(\rho\|\sigma) = D(\mathcal{N}(\rho)\|\mathcal{N}(\sigma))$ iff \exists a recovery map $\mathcal{R}_{\sigma,\mathcal{N}}^P$ such that

$$\rho = (\mathcal{R}_{\sigma,\mathcal{N}}^P \circ \mathcal{N})(\rho), \quad \sigma = (\mathcal{R}_{\sigma,\mathcal{N}}^P \circ \mathcal{N})(\sigma)$$

- This “Petz” recovery map has the following explicit form [HJPW04]:

$$\mathcal{R}_{\sigma,\mathcal{N}}^P(\omega) \equiv \sigma^{1/2} \mathcal{N}^\dagger \left((\mathcal{N}(\sigma))^{-1/2} \omega (\mathcal{N}(\sigma))^{-1/2} \right) \sigma^{1/2}$$

- Classical case: Distributions p_X and q_X and a channel $\mathcal{N}(y|x)$. Then the Petz recovery map $\mathcal{R}^P(x|y)$ is given by the Bayes theorem:

$$\mathcal{R}^P(x|y)q_Y(y) = \mathcal{N}(y|x)q_X(x)$$

where $q_Y(y) \equiv \sum_x \mathcal{N}(y|x)q_X(x)$

More on Petz recovery map

- Linear, completely positive by inspection and trace non-increasing because

$$\begin{aligned}\mathrm{Tr}\{\mathcal{R}_{\sigma,\mathcal{N}}^P(\omega)\} &= \mathrm{Tr}\{\sigma^{1/2}\mathcal{N}^\dagger\left((\mathcal{N}(\sigma))^{-1/2}\omega(\mathcal{N}(\sigma))^{-1/2}\right)\sigma^{1/2}\} \\ &= \mathrm{Tr}\{\sigma\mathcal{N}^\dagger\left((\mathcal{N}(\sigma))^{-1/2}\omega(\mathcal{N}(\sigma))^{-1/2}\right)\} \\ &= \mathrm{Tr}\{\mathcal{N}(\sigma)(\mathcal{N}(\sigma))^{-1/2}\omega(\mathcal{N}(\sigma))^{-1/2}\} \\ &\leq \mathrm{Tr}\{\omega\}\end{aligned}$$

- For $\mathcal{N}(\sigma)$ positive definite, the map perfectly recovers σ from $\mathcal{N}(\sigma)$:

$$\begin{aligned}\mathcal{R}_{\sigma,\mathcal{N}}^P(\mathcal{N}(\sigma)) &= \sigma^{1/2}\mathcal{N}^\dagger\left((\mathcal{N}(\sigma))^{-1/2}\mathcal{N}(\sigma)(\mathcal{N}(\sigma))^{-1/2}\right)\sigma^{1/2} \\ &= \sigma^{1/2}\mathcal{N}^\dagger(I)\sigma^{1/2} \\ &= \sigma\end{aligned}$$

Normalization [LW14]

For identity channel, the Petz recovery map is the identity map:

$\mathcal{R}_{\sigma, \text{id}}^P = \text{id}$. “If there’s no noise, then no need to recover”

Tensorial [LW14]

Given a tensor-product state and channel, then the Petz recovery map is a

tensor product: $\mathcal{R}_{\sigma_1 \otimes \sigma_2, \mathcal{N}_1 \otimes \mathcal{N}_2}^P = \mathcal{R}_{\sigma_1, \mathcal{N}_1}^P \otimes \mathcal{R}_{\sigma_2, \mathcal{N}_2}^P$. “Individual action suffices for ‘pretty good’ recovery of individual states”

Composition [LW14]

Given $\mathcal{N}_2 \circ \mathcal{N}_1$, then $\mathcal{R}_{\sigma, \mathcal{N}_2 \circ \mathcal{N}_1}^P = \mathcal{R}_{\sigma, \mathcal{N}_1}^P \circ \mathcal{R}_{\mathcal{N}_1(\sigma), \mathcal{N}_2}^P$. “To recover ‘pretty well’ overall, recover ‘pretty well’ from the last noise first and the first noise last”

Petz recovery map for strong subadditivity

- Strong subadditivity is a special case of monotonicity of relative entropy with $\rho = \omega_{ABC}$, $\sigma = \omega_{AC} \otimes I_B$, and $\mathcal{N} = \text{Tr}_A$
- Then $\mathcal{N}^\dagger(\cdot) = (\cdot) \otimes I_A$ and Petz recovery map is

$$\mathcal{R}_{C \rightarrow AC}^P(\tau_C) = \omega_{AC}^{1/2} \left(\omega_C^{-1/2} \tau_C \omega_C^{-1/2} \otimes I_A \right) \omega_{AC}^{1/2}$$

- Interpretation: If system A is lost but $H(B|C)_\omega = H(B|AC)_\omega$, then one can recover the full state on ABC by performing the Petz recovery map on system C of ω_{BC} , i.e.,

$$\omega_{ABC} = \mathcal{R}_{C \rightarrow AC}^P(\omega_{BC})$$

Approximate case

Approximate case would be useful for applications

Approximate case for monotonicity of relative entropy

- What can we say when $D(\rho\|\sigma) - D(\mathcal{N}(\rho)\|\mathcal{N}(\sigma)) = \varepsilon$?
- Does there exist a CPTP map \mathcal{R} that recovers σ perfectly from $\mathcal{N}(\sigma)$ while recovering ρ from $\mathcal{N}(\rho)$ approximately? [WL12]

Approximate case for strong subadditivity

- What can we say when $H(B|C)_\omega - H(B|AC)_\omega = \varepsilon$?
- Is ω_{ABC} approximately recoverable from ω_{BC} by performing a recovery map on system C alone? [WL12]

Other measures of similarity for quantum states

Trace distance

Trace distance between ρ and σ is $\|\rho - \sigma\|_1$ where $\|A\|_1 = \text{Tr}\{\sqrt{A^\dagger A}\}$.
Has a one-shot operational interpretation as the bias in success probability when distinguishing ρ and σ with an optimal quantum measurement.

Fidelity [Uhl76]

Fidelity between ρ and σ is $F(\rho, \sigma) \equiv \|\sqrt{\rho}\sqrt{\sigma}\|_1^2$. Has a one-shot operational interpretation as the probability with which a purification of ρ could pass a test for being a purification of σ .

Bures distance [Bur69]

Bures distance between ρ and σ is $D_B(\rho, \sigma) = \sqrt{2(1 - \sqrt{F(\rho, \sigma)})}$.

Breakthrough result of [FR14]

Remainder term for strong subadditivity [FR14]

\exists unitary channels \mathcal{U}_C and \mathcal{V}_{AC} such that

$$H(B|C)_\omega - H(B|AC)_\omega \geq -\log F\left(\omega_{ABC}, (\mathcal{V}_{AC} \circ \mathcal{R}_{C \rightarrow AC}^P \circ \mathcal{U}_C)(\omega_{BC})\right)$$

Nothing known from [FR14] about these unitaries! However, can conclude that $I(A; B|C)$ is small iff ω_{ABC} is approximately recoverable from system C alone after the loss of system A .

Remainder term for monotonicity of relative entropy [BLW14]

\exists unitary channels \mathcal{U} and \mathcal{V} such that

$$D(\rho||\sigma) - D(\mathcal{N}(\rho)||\mathcal{N}(\sigma)) \geq -\log F\left(\rho, (\mathcal{V} \circ \mathcal{R}_{\sigma, \mathcal{N}}^P \circ \mathcal{U})(\mathcal{N}(\rho))\right)$$

Again, nothing known from [BLW14] about \mathcal{U} and \mathcal{V} .

New result of [Wil15]

New Theorem: Let ρ and σ be such that $\text{supp}(\rho) \subseteq \text{supp}(\sigma)$ and let \mathcal{N} be a quantum channel. Then the following inequality holds

$$D(\rho \parallel \sigma) - D(\mathcal{N}(\rho) \parallel \mathcal{N}(\sigma)) \geq -\log \left[\sup_{t \in \mathbb{R}} F\left(\rho, \mathcal{R}_{\sigma, \mathcal{N}}^{P, t}(\mathcal{N}(\rho))\right) \right],$$

where $\mathcal{R}_{\sigma, \mathcal{N}}^{P, t}$ is the following rotated Petz recovery map:

$$\mathcal{R}_{\sigma, \mathcal{N}}^{P, t}(\cdot) \equiv \left(\mathcal{U}_{\sigma, t} \circ \mathcal{R}_{\sigma, \mathcal{N}}^P \circ \mathcal{U}_{\mathcal{N}(\sigma), -t} \right) (\cdot),$$

$\mathcal{R}_{\sigma, \mathcal{N}}^P$ is the Petz recovery map, and $\mathcal{U}_{\sigma, t}$ and $\mathcal{U}_{\mathcal{N}(\sigma), -t}$ are defined from $\mathcal{U}_{\omega, t}(\cdot) \equiv \omega^{it}(\cdot)\omega^{-it}$, with ω a positive semi-definite operator.

Two tools for proof: Rényi generalization of a relative entropy difference and the Hadamard three-line theorem

Rényi generalizations of a relative entropy difference

Definition from [BSW14, SBW14]

$$\tilde{\Delta}_\alpha(\rho, \sigma, \mathcal{N}) \equiv \frac{2}{\alpha'} \log \left\| \left(\mathcal{N}(\rho)^{-\alpha'/2} \mathcal{N}(\sigma)^{\alpha'/2} \otimes I_E \right) U \sigma^{-\alpha'/2} \rho^{1/2} \right\|_{2\alpha},$$

where $\alpha \in (0, 1) \cup (1, \infty)$, $\alpha' \equiv (\alpha - 1)/\alpha$, and $U_{S \rightarrow BE}$ is an isometric extension of \mathcal{N} .

Important properties

$$\lim_{\alpha \rightarrow 1} \tilde{\Delta}_\alpha(\rho, \sigma, \mathcal{N}) = D(\rho \| \sigma) - D(\mathcal{N}(\rho) \| \mathcal{N}(\sigma)).$$

$$\tilde{\Delta}_{1/2}(\rho, \sigma, \mathcal{N}) = -\log F\left(\rho, \mathcal{R}_{\sigma, \mathcal{N}}^P(\mathcal{N}(\rho))\right).$$

Hadamard three-line theorem

Let $S \equiv \{z \in \mathbb{C} : 0 \leq \operatorname{Re}\{z\} \leq 1\}$, and let $L(\mathcal{H})$ be the space of bounded linear operators acting on a Hilbert space \mathcal{H} . Let $G : S \rightarrow L(\mathcal{H})$ be a bounded map that is holomorphic on the interior of S and continuous on the boundary. Let $\theta \in (0, 1)$ and define p_θ by

$$\frac{1}{p_\theta} = \frac{1 - \theta}{p_0} + \frac{\theta}{p_1},$$

where $p_0, p_1 \in [1, \infty]$. For $k = 0, 1$ define

$$M_k = \sup_{t \in \mathbb{R}} \|G(k + it)\|_{p_k}.$$

Then

$$\|G(\theta)\|_{p_\theta} \leq M_0^{1-\theta} M_1^\theta.$$

Three (or so) line proof

Pick

$$G(z) \equiv \left([\mathcal{N}(\rho)]^{z/2} [\mathcal{N}(\sigma)]^{-z/2} \otimes I_E \right) U \sigma^{z/2} \rho^{1/2},$$
$$p_0 = 2, \quad p_1 = 1, \quad \theta \in (0, 1) \Rightarrow p_\theta = \frac{2}{1 + \theta}$$

Then

$$M_0 = \sup_{t \in \mathbb{R}} \left\| \left(\mathcal{N}(\rho)^{it/2} \mathcal{N}(\sigma)^{-it/2} \otimes I_E \right) U \sigma^{it} \rho^{1/2} \right\|_2 \leq \left\| \rho^{1/2} \right\|_2 = 1,$$

$$M_1 = \sup_{t \in \mathbb{R}} \|G(1 + it)\|_1 = \left[\sup_{t \in \mathbb{R}} F\left(\rho, \mathcal{R}_{\sigma, \mathcal{N}}^{P, t}(\mathcal{N}(\rho))\right) \right]^{1/2}.$$

Apply the three-line theorem to conclude that

$$\|G(\theta)\|_{2/(1+\theta)} \leq \left[\sup_{t \in \mathbb{R}} F\left(\rho, \mathcal{R}_{\sigma, \mathcal{N}}^{P, t}(\mathcal{N}(\rho))\right) \right]^{\theta/2}.$$

Take a negative logarithm and the limit as $\theta \searrow 0$ to conclude.

SSA refinement as a special case

Let ρ_{ABC} be a density operator acting on a finite-dimensional Hilbert space $\mathcal{H}_A \otimes \mathcal{H}_B \otimes \mathcal{H}_C$. Then the following inequality holds

$$I(A; B|C)_\rho \geq -\log \left[\sup_{t \in \mathbb{R}} F \left(\rho_{ABC}, \mathcal{R}_{C \rightarrow AC}^{P,t}(\rho_{BC}) \right) \right],$$

where $\mathcal{R}_{C \rightarrow AC}^{P,t}$ is the following rotated Petz recovery map:

$$\mathcal{R}_{C \rightarrow AC}^{P,t}(\cdot) \equiv \left(\mathcal{U}_{\rho_{AC},t} \circ \mathcal{R}_{C \rightarrow AC}^P \circ \mathcal{U}_{\rho_C,-t} \right)(\cdot),$$

the Petz recovery map $\mathcal{R}_{C \rightarrow AC}^P$ is defined as

$$\mathcal{R}_{C \rightarrow AC}^P(\cdot) \equiv \rho_{AC}^{1/2} \left[\rho_C^{-1/2}(\cdot) \rho_C^{-1/2} \otimes I_A \right] \rho_{AC}^{1/2},$$

and the partial isometric maps $\mathcal{U}_{\rho_{AC},t}$ and $\mathcal{U}_{\rho_C,-t}$ are defined as before.

Conclusions

- The result of [FR14] already had a number of important implications in quantum information theory.
- The new result in [Wil15] applies to relative entropy differences, has a brief proof, and improves our understanding of the input and output unitaries (but see [SFR15] for the special case of SSA)
- By building on [SFR15, Wil15], we can now generalize these results: there is a universal recovery map which depends only on σ and \mathcal{N} and has the form [SRWW15]:

$$X \rightarrow \int \mu(dt) \mathcal{R}_{\sigma, \mathcal{N}}^{P,t}(X)$$

for some probability measure μ .

- It is still conjectured that the recovery map can be the Petz recovery map alone (not a rotated Petz map).

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