

# Quantum polar codes for arbitrary channels

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# The Quantum Coding Problem

We have some idea of good rates for quantum communication over quantum channels  
*(and in some cases, we know capacity)*

**Quantum turbo codes** and **quantum LDPC codes** are attempts at explicit constructions, but it seems difficult to prove that they are capacity-achieving.

**Polar codes** are a promising code construction in the classical world, so why not explore their quantum generalization?

Result is a **near-explicit, capacity-achieving scheme** for some quantum channels

# Overview

## Quantum codes from classical-quantum codes!

Using Arikan's code construction (encoder/decoder) directly would have suboptimal performance for general quantum channels...

Though, a generalization of the decoder to classical input, quantum output channels is helpful for us

First discuss polar codes for classical-quantum channels

Then discuss quantum polar codes from classical-quantum ones

Finally outline how quantum polar codes exhibit superactivation

# Classical-Quantum Channels

Begin with a binary-input, classical-quantum channel:

$$W : x \rightarrow \rho_x$$

One channel parameter is **symmetric Holevo information**:

$$\begin{aligned} I(W) &\equiv I(X; B) \\ &= H((\rho_0 + \rho_1)/2) - H(\rho_0)/2 - H(\rho_1)/2 \end{aligned}$$

Evaluate  $I(X;B)$  with respect to

$$\frac{1}{2} (|0\rangle\langle 0|^X \otimes \rho_0^B + |1\rangle\langle 1|^X \otimes \rho_1^B)$$

Equal to one for *perfect channels* and zero for *useless channels*

# Fidelity Channel Parameter

**Fidelity** characterizes **distinguishability** of two output states:

$$\begin{aligned} F(W) &\equiv F(\rho_0, \rho_1) \\ &= \|\sqrt{\rho_0}\sqrt{\rho_1}\|_1^2 \end{aligned}$$

$F(W) = 0$  if states are *perfectly distinguishable*

$F(W) = 1$  if states are *not distinguishable*

Generalizes classical fidelity (Bhattacharya parameter)

# Quantum Hypothesis Testing

Suppose someone else prepares  $\rho$  or  $\sigma$  with equal probability and then hands you the prepared system

**Your task:**

Perform a measurement to determine which state was prepared

Your measurement will be of the following form:

$$\{\Lambda_\rho, \Lambda_\sigma\} \quad \text{where} \quad \Lambda_\rho + \Lambda_\sigma = I$$

Error probability is then

$$p_{\text{err}} = \frac{1}{2} (\text{Tr}\{\Lambda_\rho \sigma\} + \text{Tr}\{\Lambda_\sigma \rho\})$$

Fidelity serves as an **upper bound** on error probability in a **quantum hypothesis test** that attempts to distinguish  $\rho_0$  from  $\rho_1$ :

$$P_e \leq \frac{\sqrt{F(W)}}{2}$$

# Polar Coding Scheme

Encoding circuit—same as Arikan's, though use fidelity for polar coding rule

Send information bits through the good channels

Send frozen (ancilla) bits through the bad channels

## Quantum Successive Cancellation Decoder

performs quantum hypothesis tests  
to make decisions on the information bits

**Key tool** in the proof that this scheme works  
is Pranab Sen's “**non-commutative union bound**”:

$$1 - \text{Tr}\{\Pi_N \cdots \Pi_1 \rho \Pi_1 \cdots \Pi_N\} \leq 2 \sqrt{\sum_{i=1}^N \text{Tr}\{(I - \Pi_i)\rho\}}$$

This leads to a near-explicit capacity-achieving scheme

# Towards Quantum Polar Codes

Use amplitude and phase coding ideas of Renes and Boileau

Suppose we're trying to code for a channel  $\mathcal{N}^{A' \rightarrow B}$

Consider coding for a cq channel:

$$W_A : z \rightarrow \mathcal{N}^{A' \rightarrow B} (|z\rangle \langle z|)$$

Suppose there is a good decoding POVM  $\{\Lambda_z\}$

$$\text{Tr}\{\Lambda_z \mathcal{N}(|z\rangle \langle z|)\} \geq 1 - \epsilon$$



# Towards Quantum Polar Codes

Similarly, consider the cq phase channel with quantum side info.

$$W_P : x \rightarrow (Z^x)^C U_{\mathcal{N}}^{A' \rightarrow BE} |\Phi\rangle^{CA'}$$

where the state on  $CA'$  is the maximally entangled Bell state (just assuming for the moment that Bob has  $C$ )

Suppose there is a good decoding POVM  $\{\Gamma_x\}$

$$\text{Tr}\{\Gamma_x \sigma_x^{BC}\} \geq 1 - \epsilon$$

# First Decoding Step

Initial state: 
$$\sum_z \sqrt{p(z)} |z\rangle^A |z\rangle^{A'}$$

State sent through channel: 
$$\sum_z \sqrt{p(z)} |z\rangle^A |\psi_z\rangle^{BE}$$

Coherent version of cq decoder: 
$$\sum_z \sqrt{\Lambda_z}^B \otimes |z\rangle^C$$

State after 1<sup>st</sup> decoding step:

$$\approx \sum_z \sqrt{p(z)} |z\rangle^A |\psi_z\rangle^{BE} |z\rangle^C$$

This approximation follows from performance of cq code

# Second Decoding Step

Rewrite the state using the Fourier basis:

$$\begin{aligned}\sum_z \sqrt{p(z)} |z\rangle^A |\psi_z\rangle^{BE} |z\rangle^C &= \sum_x |\tilde{x}\rangle \langle \tilde{x}|^A \sum_z \sqrt{p(z)} |z\rangle^A |\psi_z\rangle^{BE} |z\rangle^C \\ &= \sum_x |\tilde{x}\rangle^A \sum_z \langle \tilde{x}|z\rangle \sqrt{p(z)} |\psi_z\rangle^{BE} |z\rangle^C \\ &= \frac{1}{\sqrt{d}} \sum_x |\tilde{x}\rangle^A (Z^{-x})^C \sum_z \sqrt{p(z)} |\psi_z\rangle^{BE} |z\rangle^C\end{aligned}$$

Coherent version of cq phase decoder:  $\sum_x \sqrt{\Gamma_x}^{BC} \otimes |\tilde{x}\rangle^D$

State after 2<sup>nd</sup> decoding step:

$$\approx \frac{1}{\sqrt{d}} \sum_x |\tilde{x}\rangle^A (Z^{-x})^C \sum_z \sqrt{p(z)} |\psi_z\rangle^{BE} |z\rangle^C |\tilde{x}\rangle^D$$

# Final Decoding Step

State after 2<sup>nd</sup> decoding step:

$$\approx \frac{1}{\sqrt{d}} \sum_x |\tilde{x}\rangle^A (Z^{-x})^C \sum_z \sqrt{p(z)} |\psi_z\rangle^{BE} |z\rangle^C |\tilde{x}\rangle^D$$

Bob performs controlled-phase gates on C and D:

$$\frac{1}{\sqrt{d}} \sum_x |\tilde{x}\rangle^A \sum_z \sqrt{p(z)} |\psi_z\rangle^{BE} |z\rangle^C |\tilde{x}\rangle^D$$

Rewrite this state as

$$\frac{1}{\sqrt{d}} \sum_x |\tilde{x}\rangle^A |\tilde{x}\rangle^D \otimes \sum_z \sqrt{p(z)} |\psi_z\rangle^{BE} |z\rangle^C$$

**Decoding successful!** (b/c maximal ent. shared between A and D)

# Quantum Polar Codes

Use amplitude and phase encoding ideas of Renes

Build quantum polar codes from cq channels:

$$W_A : z \rightarrow \mathcal{N}^{A' \rightarrow B} (|z\rangle \langle z|)$$

$$W_P : x \rightarrow (Z^x)^C U_{\mathcal{N}}^{A' \rightarrow BE} |\Phi\rangle^{CA'}$$

Good for Amp, bad for Phase: send  $|+\rangle$  states into these

Good for Amp, good for Phase: send information qubits into these

Bad for Amp, bad for Phase: send halves of ebits into these

Bad for Amp, good for Phase: send ancilla qubits  $|0\rangle$  into these

# Construction (Ctd.)

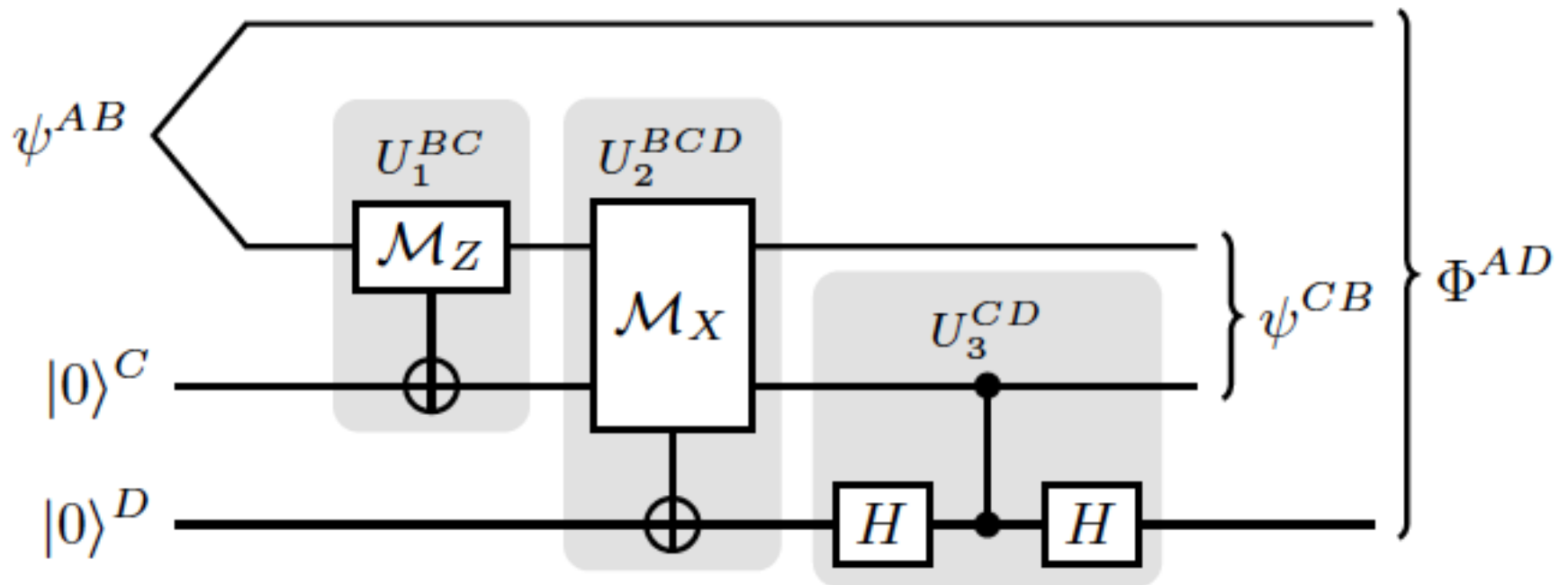
$N \cdot I(Z;B)$  channels **good for Amplitude**

$N \cdot I(X;BC)$  channels **good for Phase**

Can show that **net rate** of quantum communication is

$$I(Z;B) + I(X;BC) - 1 = I(A \rangle B)$$

Decoder operates coherently



# Superactivation with near-explicit codes

Get near explicit codes for **superactivation** as a bonus!

Example of channels found by Smith and Yard  
each have **4-dimensional inputs**,

Giving a **16-dimensional** input space for joint channel

Factor this as a tensor product of **4 qubit input spaces**,  
and then apply a slightly modified version  
of the amplitude and phase construction

Coherently decode the amplitude and phase variables in the order:

Z1,  
Z2 | Z1,  
Z3 | Z1 Z2,  
Z4 | Z1 Z2 Z3,  
X1 | Z1 Z2 Z3 Z4,  
X2 | Z1 Z2 Z3 Z4 X1,  
X3 | Z1 Z2 Z3 Z4 X1 X2,  
X4 | Z1 Z2 Z3 Z4 X1 X2 X3,

# Conclusion

**Polar coding** gives a near-explicit, capacity-achieving scheme for quantum communication

Even gives a near-explicit scheme for **superactivation**

*Most important open problem:*

Show how to make the decoder **efficient**  
(progress in Renes *et al.* arXiv:1109.3195 for Pauli channels)

*Other important problems:*

- 1) Which channels are the good ones?
- 2) Extend to other scenarios